Normal forms and symmetries of (2,5) distributions: 100 years after Cartan

> Michail Zhitomirskii Technion

January 30, 2011 Hiroshima

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

Congratulations to Professor Keizo Yamaguchi Professor Reiko Miyaoka

▲□▶ ▲□▶ ▲ 臣▶ ★ 臣▶ 三臣 - のへぐ

 E. Cartan, Les systemes de Pfaff a cinque variables et les equations aux derivees partielles du second ordre, Ann. Sci. Ecole Normale, 1910

▲□▶ ▲□▶ ▲ 臣▶ ★ 臣▶ 三臣 - のへぐ

- E. Cartan, Les systemes de Pfaff a cinque variables et les equations aux derivees partielles du second ordre, Ann. Sci. Ecole Normale, 1910
- Topic (in terms of vector fields, Cartan used equivalent language of 1-forms):

2-planes in $T\mathbb{R}^5 = (2,5)$ distributions.

- E. Cartan, Les systemes de Pfaff a cinque variables et les equations aux derivees partielles du second ordre, Ann. Sci. Ecole Normale, 1910
- Topic (in terms of vector fields, Cartan used equivalent language of 1-forms):

2-planes in $T\mathbb{R}^5 = (2,5)$ distributions.

► Most famous results concern (2, 3, 5) distributions: span(V₁, V₂) such that the vectors

 $V_1, V_2, [V_1, V_2], [V_1, [V_1, V_2]], [V_2, [V_1, V_2]]$

are linearly independent at any point (generic growth vector).



SELECTED CARTAN'S RESULTS on symmetries of (2,3,5) distributions (germs at $0 \in \mathbb{R}^5$)



SELECTED CARTAN'S RESULTS on symmetries of (2,3,5) distributions (germs at $0 \in \mathbb{R}^5$)

- ► The maximal possible dimension of the group of local symmetries is 14. All distributions with 14-dim symmetry group are diffeomorphic; the symmetry group is simple (it is G₂).
- Terminology (not Cartan's): such distributions are called flat (2, 3, 5) distributions.

SELECTED CARTAN's RESULTS on symmetries of (2,3,5) distributions (germs at $0 \in \mathbb{R}^5$)

► The maximal possible dimension of the group of local symmetries is 14. All distributions with 14-dim symmetry group are diffeomorphic; the symmetry group is simple (it is G₂).

►

- Terminology (not Cartan's): such distributions are called flat (2, 3, 5) distributions.
- ► A distribution D is flat if and only if it admits a nilpotent (2,3,5) basis V₁, V₂:

 $D = span(V_1, V_2)$: all length ≥ 4 Lie brackets of V_1, V_2 are equal to 0.

 The maximal possible dimension of the symmetry group of a non-flat distribution is 7.

▲□▶ ▲圖▶ ▲≣▶ ▲≣▶ = = -の�?

- The maximal possible dimension of the symmetry group of a non-flat distribution is 7.
- In the local classification of distributions with 7-dimensional or 6-dimensional group of symmetries there is exactly one modulus.

 Cartan proves these and other results using his famous Cartan tensor.

- Cartan proves these and other results using his famous Cartan tensor.
- Given a (2, 3, 5) distribution D Cartan constructs a map, by a long procedure: Cartan's method (way, road, approach?)

$$\mathbb{R}^5
i p
ightarrow C_D^{(4)}(p;x_1,x_2)$$
 (Cartan's tensor at p)

- Cartan proves these and other results using his famous Cartan tensor.
- Given a (2, 3, 5) distribution D Cartan constructs a map, by a long procedure: Cartan's method (way, road, approach?)

$$\mathbb{R}^5
i p
ightarrow C_D^{(4)}(p;x_1,x_2)$$
 (Cartan's tensor at p)

if the germ of D at p is diffeomorphic to the germ of D at p
 is diffeomorphic to the germ of D
 at p
 then C_D⁽⁴⁾(p; x₁, x₂) and C_D⁽⁴⁾(p
 ; x₁, x₂) are linearly equivalent;

- Cartan proves these and other results using his famous Cartan tensor.
- Given a (2,3,5) distribution D Cartan constructs a map, by a long procedure: Cartan's method (way, road, approach?)

$$\mathbb{R}^5
i p
ightarrow C_D^{(4)}(p;x_1,x_2)$$
 (Cartan's tensor at p)

- ▶ if the germ of D at p is diffeomorphic to the germ of D at p then C_D⁽⁴⁾(p; x₁, x₂) and C_D⁽⁴⁾(p̃; x₁, x₂) are linearly equivalent;
- ► ⇒ the Cartan tensor C_D⁽⁴⁾(0; x₁, x₂) defined modulo linear equivalence is the Cartan invariant of the distribution germ at 0 ∈ ℝ⁵ (wrt local diffeomorphisms preserving 0).

- Cartan proves these and other results using his famous Cartan tensor.
- Given a (2, 3, 5) distribution D Cartan constructs a map, by a long procedure: Cartan's method (way, road, approach?)

$$\mathbb{R}^5
i p
ightarrow C_D^{(4)}(p;x_1,x_2)$$
 (Cartan's tensor at p)

- if the germ of D at p is diffeomorphic to the germ of D at p
 is diffeomorphic to the germ of D
 at p
 then C_D⁽⁴⁾(p; x₁, x₂) and C_D⁽⁴⁾(p
 ; x₁, x₂) are linearly equivalent;
- ▶ ⇒ the Cartan tensor $C_D^{(4)}(0; x_1, x_2)$ defined modulo linear equivalence is the Cartan invariant of the distribution germ at $0 \in \mathbb{R}^5$ (wrt local diffeomorphisms preserving 0).
- ▶ if *D* is the germ at $0 \in \mathbb{R}^5$ then *D* is flat if and only $C_D^{(4)}(p; x_1, x_2) = 0$ for all *p* close to 0.

 I will show how the mentioned Cartan's results and a few new results can be obtained and easily explained using exact normal forms for (2,3,5) distributions,

▲□▶ ▲□▶ ▲ 臣▶ ★ 臣▶ 三臣 - のへぐ

- ► I will show how the mentioned Cartan's results and a few new results can be obtained and easily explained using exact normal forms for (2, 3, 5) distributions,
- and how exact normal forms can be constructed using the "parallel" normalization procedure along with quasi-homogeneous filtration in the space of infinite jets.

- I will show how the mentioned Cartan's results and a few new results can be obtained and easily explained using exact normal forms for (2,3,5) distributions,
- and how exact normal forms can be constructed using the "parallel" normalization procedure along with quasi-homogeneous filtration in the space of infinite jets.
- Instead of usual jets one should work with quasi-jets wrt natural weights

$$w(x_1) = w(x_2) = 1$$
, $w(x_3) = 2$, $w(x_4) = w(x_5) = 3$

- I will show how the mentioned Cartan's results and a few new results can be obtained and easily explained using exact normal forms for (2,3,5) distributions,
- and how exact normal forms can be constructed using the "parallel" normalization procedure along with quasi-homogeneous filtration in the space of infinite jets.
- Instead of usual jets one should work with quasi-jets wrt natural weights

$$w(x_1) = w(x_2) = 1$$
, $w(x_3) = 2$, $w(x_4) = w(x_5) = 3$

otherwise much longer computation and weaker results.

- ► I will show how the mentioned Cartan's results and a few new results can be obtained and easily explained using exact normal forms for (2, 3, 5) distributions,
- and how exact normal forms can be constructed using the "parallel" normalization procedure along with quasi-homogeneous filtration in the space of infinite jets.
- Instead of usual jets one should work with quasi-jets wrt natural weights

$$w(x_1) = w(x_2) = 1$$
, $w(x_3) = 2$, $w(x_4) = w(x_5) = 3$

- otherwise much longer computation and weaker results.
- The normalization procedure, leading to exact normal form, consists of two steps (if one uses the usual filtration: 7 or 8 steps).

 The most important is the first step. Though it gives only "almost exact" normal form (exact up to a finite dimensional group of transformations),

it gives a simple explanation of Cartan invariant and allows to generalize it.

The most important is the first step. Though it gives only "almost exact" normal form (exact up to a finite dimensional group of transformations),

it gives a simple explanation of Cartan invariant and allows to generalize it.

- Modulo quasi-homogeneity, the first step of the normalization procedure is not more than the classical "normalization by the principal part", for example the resonant normal form serving for all germs of singular vector field germs with a fixed linear approximation at 0 (Poincare, Dulac).
- ► The role of linear approximation: the nilpotent approximation of a (2,3,5) distribution D which is the symbol of D (nilpotent graded (2,3,5) Lie algebra) expressed in terms of quasi-homogeneous degree -1 vector fields.

► The first step of the normalization procedure based on quasi-homogeneity is, to some extend, "diffeomorphic" to Tanaka prolongation of the nilpotent (2, 3, 5) algebra, and probably some of results I will tell about can be obtained developing Tanaka prolongation, but in my opinion developing the method going back to Poincare and Dulak is simpler.

- ► The first step of the normalization procedure based on quasi-homogeneity is, to some extend, "diffeomorphic" to Tanaka prolongation of the nilpotent (2, 3, 5) algebra, and probably some of results I will tell about can be obtained developing Tanaka prolongation, but in my opinion developing the method going back to Poincare and Dulak is simpler.
- The generalized Cartan invariant allows to analyze the Lie algebras of all possible groups of symmetries preserving 0 (a very important subgroup of the whole symmetry group) and to prove

- ► The first step of the normalization procedure based on quasi-homogeneity is, to some extend, "diffeomorphic" to Tanaka prolongation of the nilpotent (2, 3, 5) algebra, and probably some of results I will tell about can be obtained developing Tanaka prolongation, but in my opinion developing the method going back to Poincare and Dulak is simpler.
- The generalized Cartan invariant allows to analyze the Lie algebras of all possible groups of symmetries preserving 0 (a very important subgroup of the whole symmetry group) and to prove
- ► Theorem. As an abstract Lie algebra, the algebra of infinitesimal symmetries vanishing at 0, of any non-flat (2,3,5) distribution, is either {0}, or 1-dimensional, or 2-dimensional non-Abelian, or the 3-dimensional algebra sl₂(ℝ) (the last case seems to be new).

Continuation of the theorem. As a Lie algebra of singular vector fields, in each of these cases it is linearizable, i.e. in suitable local coordinates it consists of linear vector fields, even though each of them is resonant.

Any vanishing at 0 infinitesinmal symmetry of any non-flat (2,3,5) distribution has, in suitable coordinates, the form

$$V_{A}: \begin{pmatrix} \dot{x}_{1} \\ \dot{x}_{2} \end{pmatrix} = A \begin{pmatrix} x_{1} \\ x_{2} \end{pmatrix}, \quad \dot{x}_{3} = (traceA)x_{3},$$
$$\begin{pmatrix} \dot{x}_{4} \\ \dot{x}_{5} \end{pmatrix} = (A + traceA \cdot I) \begin{pmatrix} x_{4} \\ \dot{x}_{5} \end{pmatrix}$$

where a constant 2×2 matrix A is one of the follows:

$$\begin{array}{ccc} \bullet & \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, & \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}: & \text{vector fields } b_{1:1}^{\pm} \\ & \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}: & \text{vector field } b_{0:1} \\ & \begin{pmatrix} p & 0 \\ 0 & -q \end{pmatrix}: & 1 \leq p < q: & \text{vector field } b_{p:q} \\ & \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}: & \text{vector field } b_{0:0} \end{array}$$

Either one of the eigenvalues is 0 or their ratio is a negative rational number.

・ロト ・ 日 ・ モ ト ・ モ ・ うへぐ

$$b_{1:1}^{+} = x_2 \frac{\partial}{\partial x_1} - x_1 \frac{\partial}{\partial x_2} + x_5 \frac{\partial}{\partial x_4} - x_4 \frac{\partial}{\partial x_5}$$

$$b_{1:1}^{-} = x_1 \frac{\partial}{\partial x_1} - x_2 \frac{\partial}{\partial x_2} + x_4 \frac{\partial}{\partial x_4} - x_5 \frac{\partial}{\partial x_5}$$

$$b_{1:0} = x_2 \frac{\partial}{\partial x_2} + x_3 \frac{\partial}{\partial x_3} + x_4 \frac{\partial}{\partial x_4} + 2x_5 \frac{\partial}{\partial x_5}$$

$$b_{p:q} = p x_1 \frac{\partial}{\partial x_1} - q x_2 \frac{\partial}{\partial x_2} + (p-q) x_3 \frac{\partial}{\partial x_3} + (2p-q) x_4 \frac{\partial}{\partial x_4} + (q-2p) x_5 \frac{\partial}{\partial x_5}$$

$$b_{0:0} = x_1 \frac{\partial}{\partial x_2} + x_4 \frac{\partial}{\partial x_5} .$$

▲□▶ ▲□▶ ▲三▶ ▲三▶ 三三 - のへで

►
$$b_{1:1}^+ = x_2 \frac{\partial}{\partial x_1} - x_1 \frac{\partial}{\partial x_2} + x_5 \frac{\partial}{\partial x_4} - x_4 \frac{\partial}{\partial x_5}$$

 $b_{1:1}^- = x_1 \frac{\partial}{\partial x_1} - x_2 \frac{\partial}{\partial x_2} + x_4 \frac{\partial}{\partial x_4} - x_5 \frac{\partial}{\partial x_5}$
 $b_{1:0} = x_2 \frac{\partial}{\partial x_2} + x_3 \frac{\partial}{\partial x_3} + x_4 \frac{\partial}{\partial x_4} + 2x_5 \frac{\partial}{\partial x_5}$
 $b_{p:q} = px_1 \frac{\partial}{\partial x_1} - qx_2 \frac{\partial}{\partial x_2} + (p-q)x_3 \frac{\partial}{\partial x_3} + (2p-q)x_4 \frac{\partial}{\partial x_4} + (q-2p)x_5 \frac{\partial}{\partial x_5}$
 $b_{0:0} = x_1 \frac{\partial}{\partial x_2} + x_4 \frac{\partial}{\partial x_5}.$

 Lie algebras of infinitesimal symmetries vanishing at 0 of any non-flat distribution:

- 1-dim (any of these vector field)
- 2-dim non-Abelian $span(b_{1:0}, b_{0:0})$
- $sl_2(\mathbb{R})$: $span(b_{1:1}^+, b_{1:1}^-, b_{0:0}) =$
- = all linear traceless vector fields.

Additional important statement:

The almost exact normal form and the exact normal form are parameterized by a function $C(x_1, ..., x_5)$ in a certain ideal in the ring of function germs. They hold in the same coordinates as the coordinates in which all symmetries vanishing at 0 are linear. These symmetries annihilate $C(x_1, ..., x_5)$.

The information about the Lie algebra of infinitesimal symmetries vanishing at 0, up to diffeomorphisms rather than isomorphisms, allows to classify (easily) all possible complete symmetry algebras (including infinitesimal symmetries not vanishing at 0), and for each of them to classify distributions with this symmetry algebra.

 EXAMLE: from one symmetry to 6-dim or 7-dim algebra of symmetries

 EXAMLE: from one symmetry to 6-dim or 7-dim algebra of symmetries

► Assume we have an infinitesimal symmetry $b_{1:0} = x_2 \frac{\partial}{\partial x_2} + x_3 \frac{\partial}{\partial x_3} + x_4 \frac{\partial}{\partial x_4} + 2x_5 \frac{\partial}{\partial x_5}$

- EXAMLE: from one symmetry to 6-dim or 7-dim algebra of symmetries
- ► Assume we have an infinitesimal symmetry $b_{1:0} = x_2 \frac{\partial}{\partial x_2} + x_3 \frac{\partial}{\partial x_3} + x_4 \frac{\partial}{\partial x_4} + 2x_5 \frac{\partial}{\partial x_5}$
- ► The function C(x) in the almost exact normal form starts with quasi-degree ≥ 4 and satisfies b_{1:0}(C(x) = 0. Consequently

$$C = C(x_1) = \pm x_1^m + r_1 x_1^{m+1} + r_2 x_1^{m+2} + h.o.t.(quasi)$$
 $m \ge 4$

- EXAMLE: from one symmetry to 6-dim or 7-dim algebra of symmetries
- ► Assume we have an infinitesimal symmetry $b_{1:0} = x_2 \frac{\partial}{\partial x_2} + x_3 \frac{\partial}{\partial x_3} + x_4 \frac{\partial}{\partial x_4} + 2x_5 \frac{\partial}{\partial x_5}$
- The function C(x) in the almost exact normal form starts with quasi-degree ≥ 4 and satisfies b_{1:0}(C(x) = 0. Consequently C = C(x₁) = ±x₁^m + r₁x₁^{m+1} + r₂x₁^{m+2} + h.o.t.(quasi) m ≥ 4
- ► The second step normalization gives exact normal form: the same with r₁ = 0:

$$C = C(x_1) = \pm x_1^m + rx_1^{m+2} + h.o.t(quasi), \quad m \ge 4$$

- EXAMLE: from one symmetry to 6-dim or 7-dim algebra of symmetries
- ► Assume we have an infinitesimal symmetry $b_{1:0} = x_2 \frac{\partial}{\partial x_2} + x_3 \frac{\partial}{\partial x_3} + x_4 \frac{\partial}{\partial x_4} + 2x_5 \frac{\partial}{\partial x_5}$
- The function C(x) in the almost exact normal form starts with quasi-degree ≥ 4 and satisfies b_{1:0}(C(x) = 0. Consequently C = C(x₁) = ±x₁^m + r₁x₁^{m+1} + r₂x₁^{m+2} + h.o.t.(quasi) m ≥ 4
- ► The second step normalization gives exact normal form: the same with r₁ = 0:

$$C = C(x_1) = \pm x_1^m + rx_1^{m+2} + h.o.t(quasi), m \ge 4$$

 One gets immediately, from the exact normal form, the infinitesimal symmetry

$$b_{0:0} = x_1 \frac{\partial}{\partial x_2} + x_4 \frac{\partial}{\partial x_5}$$

- EXAMLE: from one symmetry to 6-dim or 7-dim algebra of symmetries
- ► Assume we have an infinitesimal symmetry $b_{1:0} = x_2 \frac{\partial}{\partial x_2} + x_3 \frac{\partial}{\partial x_3} + x_4 \frac{\partial}{\partial x_4} + 2x_5 \frac{\partial}{\partial x_5}$
- The function C(x) in the almost exact normal form starts with quasi-degree ≥ 4 and satisfies b_{1:0}(C(x) = 0. Consequently C = C(x₁) = ±x₁^m + r₁x₁^{m+1} + r₂x₁^{m+2} + h.o.t.(quasi) m ≥ 4
- ► The second step normalization gives exact normal form: the same with r₁ = 0:

$$C = C(x_1) = \pm x_1^m + rx_1^{m+2} + h.o.t(quasi), m \ge 4$$

 One gets immediately, from the exact normal form, the infinitesimal symmetry

$$b_{0:0} = x_1 \frac{\partial}{\partial x_2} + x_4 \frac{\partial}{\partial x_5}$$

and 4 linearly independent infinitesimal symmetries which do not vanish at 0.

•
$$C = C(x_1) = \pm x_1^m + rx_1^{m+2} + h.o.t(quasi), m \ge 4$$

▲□▶ ▲□▶ ▲三▶ ▲三▶ 三三 - のへで

•
$$C = C(x_1) = \pm x_1^m + rx_1^{m+2} + h.o.t(quasi), m \ge 4$$

 All together we have at least 6-dim algebra of infinitesimal symmetries. For a non-flat distribution we cannot have one more infinitesimal symmetry vanishing at 0.

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□

•
$$C = C(x_1) = \pm x_1^m + rx_1^{m+2} + h.o.t(quasi), m \ge 4$$

- All together we have at least 6-dim algebra of infinitesimal symmetries. For a non-flat distribution we cannot have one more infinitesimal symmetry vanishing at 0.
- Therefore for non-flat distributions b_{1:0} "gives birth" to 5 more infinitesimal symmetries, and the whole symmetry algebra has dimension 6 or 7.

◆□▶ ◆□▶ ◆ □▶ ◆ □▶ - □ - のへぐ

►
$$C = C(x_1) = \pm x_1^m + rx_1^{m+2} + h.o.t(quasi), m \ge 4$$

- All together we have at least 6-dim algebra of infinitesimal symmetries. For a non-flat distribution we cannot have one more infinitesimal symmetry vanishing at 0.
- Therefore for non-flat distributions b_{1:0} "gives birth" to 5 more infinitesimal symmetries, and the whole symmetry algebra has dimension 6 or 7.
- It has dimension 7 iff the distribution is homogeneous (the symmetry group acts transitively).

•
$$C = C(x_1) = \pm x_1^m + rx_1^{m+2} + h.o.t(quasi), m \ge 4$$

- All together we have at least 6-dim algebra of infinitesimal symmetries. For a non-flat distribution we cannot have one more infinitesimal symmetry vanishing at 0.
- Therefore for non-flat distributions b_{1:0} "gives birth" to 5 more infinitesimal symmetries, and the whole symmetry algebra has dimension 6 or 7.
- It has dimension 7 iff the distribution is homogeneous (the symmetry group acts transitively).
- It is so iff m = 4:

$$C = C(x_1) = \pm x_1^4 + rx_1^6 + h.o.t(qiuasi), m \ge 4$$

r can be arbitrary and h.o.t. are uniquely determined by r.

From the data we have we obtain that if the distribution is homogeneous and non-flat then a part of structure equations in the 7-dim algebra is as follows:

$$[a_1, a_2] = a_3, \ [a_1, a_3] = a_4, \ [a_2, a_3] = a_5$$

 $[a_1, b_{1:0}] = 0, \ [a_1, b_{0:0}] = a_2, \ [a_2, b_{1:0}] = a_2, \ [a_2, b_{0:0}] = 0$
 $[b_{1:0}, b_{0:0}] = b_{0:0}$

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□

From the data we have we obtain that if the distribution is homogeneous and non-flat then a part of structure equations in the 7-dim algebra is as follows:

$$[a_1, a_2] = a_3, \ [a_1, a_3] = a_4, \ [a_2, a_3] = a_5$$

 $[a_1, b_{1:0}] = 0, \ [a_1, b_{0:0}] = a_2, \ [a_2, b_{1:0}] = a_2, \ [a_2, b_{0:0}] = 0$
 $[b_{1:0}, b_{0:0}] = b_{0:0}$

and (as an exercise on Jacobi identity) we obtain that the algebra is isomorphic to one with the structure equations above and the structure equations:

$$[a_1, a_4] = Ba_3 - Cb_{0:0}, \quad [a_1, a_5] = [a_2, a_4] = 0, \quad [a_2, a_5] = 0.$$

(remaining structure equations follow from Jacobi identity).

•
$$C = C(x_1) = \pm x_1^4 + rx_1^6 + h.o.t, m \ge 4$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ ◆□▶

•
$$C = C(x_1) = \pm x_1^4 + rx_1^6 + h.o.t, m \ge 4$$

▶ It is a solvable algebra with "decrease vector" (7,5,1,0).

•
$$C = C(x_1) = \pm x_1^4 + rx_1^6 + h.o.t, m \ge 4$$

- ▶ It is a solvable algebra with "decrease vector" (7,5,1,0).
- ▶ The couple (*B*, *C*) is an invariant modulo scaling

$$B \to k^2 B, \ C \to k^4 C$$

therefore the sign of B and the number

$$\lambda = B^2/C$$

is an invariant. It is uniquely determined by r.

• r takes any value, and λ takes any value except

$$\lambda \neq \frac{100}{9}.$$

•
$$C = C(x_1) = \pm x_1^4 + rx_1^6 + h.o.t, m \ge 4$$

- ▶ It is a solvable algebra with "decrease vector" (7,5,1,0).
- ▶ The couple (B, C) is an invariant modulo scaling

$$B \to k^2 B, \ C \to k^4 C$$

therefore the sign of B and the number

$$\lambda = B^2/C$$

is an invariant. It is uniquely determined by r.

• r takes any value, and λ takes any value except

$$\lambda \neq \frac{100}{9}.$$

▶ \implies the (2,3,5) distribution with the above 7-dim symmetry algebra and $\lambda = 100/9$ is flat.

 Corresponding Cartan's results: Chapter IX. It takes 6 pages to prove:

- Corresponding Cartan's results: Chapter IX. It takes 6 pages to prove:
- If Cartan tensor is linearly equivalent to x₁⁴ at any point then all Cartan's invariants fundamentaux are defined by one function *I* = *I*(x₁) (Cartan's notation);

- Corresponding Cartan's results: Chapter IX. It takes 6 pages to prove:
- If Cartan tensor is linearly equivalent to x₁⁴ at any point then all Cartan's invariants fundamentaux are defined by one function *I* = *I*(x₁) (Cartan's notation);

► the symmetry group is 6-dimensional if $I \neq const$ and 7-dimensional if $I \equiv const$

- Corresponding Cartan's results: Chapter IX. It takes 6 pages to prove:
- ► If Cartan tensor is linearly equivalent to x₁⁴ at any point then all Cartan's invariants fundamentaux are defined by one function *I* = *I*(x₁) (Cartan's notation);
- ► the symmetry group is 6-dimensional if *I* ≠ const and 7-dimensional if *I* ≡ const
- the normal form:

$$\omega_1 = dx_1 + \frac{7}{3}Ix_3dx_2 + x_4dx_3 - \left(\frac{1}{2}x_4^2 + \frac{2}{3}Ix_3^2 - \frac{1}{2}(1 + I^2 - I'')x_2^2\right)dx_5,$$

$$\omega_2 = dx_2 - x_3dx_5, \quad \omega_3 = -dx_3 + x_4dx_5.$$

 EXAMPLE: homogeneous non-flat distributions with the symmetry

$$b_{1:1}^+ = x_2 \frac{\partial}{\partial x_1} - x_1 \frac{\partial}{\partial x_2} + x_5 \frac{\partial}{\partial x_4} - x_4 \frac{\partial}{\partial x_5}$$

 EXAMPLE: homogeneous non-flat distributions with the symmetry

$$b_{1:1}^+ = x_2 \frac{\partial}{\partial x_1} - x_1 \frac{\partial}{\partial x_2} + x_5 \frac{\partial}{\partial x_4} - x_4 \frac{\partial}{\partial x_5}$$

► The function C(x) in the almost exact normal form satisfies $b^+_{1:1}(C(x)) = 0$

and starts with terms of quasi-degree \geq 4. It follows

$$C(x) = \\ \pm (x_1^2 + x_2^2)^m + r_1(x_1^2 + x_2^2)^{m+1} + r_2x_3(x_1^2 + x_2^2)^2 + h.o.t(quasi)$$

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□

 EXAMPLE: homogeneous non-flat distributions with the symmetry

$$b_{1:1}^+ = x_2 \frac{\partial}{\partial x_1} - x_1 \frac{\partial}{\partial x_2} + x_5 \frac{\partial}{\partial x_4} - x_4 \frac{\partial}{\partial x_5}$$

► The function C(x) in the almost exact normal form satisfies $b^+_{1:1}(C(x)) = 0$

and starts with terms of quasi-degree \geq 4. It follows

$$C(x) = \\ \pm (x_1^2 + x_2^2)^m + r_1(x_1^2 + x_2^2)^{m+1} + r_2x_3(x_1^2 + x_2^2)^2 + h.o.t(quasi)$$

▲□▶ ▲□▶ ▲三▶ ▲三▶ - 三 - のへで

The second step normalization gives exact normal form: the same with r₂ = 0:

$$C(x) = \pm (x_1^2 + x_2^2)^m + r(x_1^2 + x_2^2)^{m+1} + h.o.t.(quasi)$$

The theorem on the all possible Lie algebras of vanishing at 0 infinitesimal symmetries implies that b⁺_{1:1} is the only such symmetry. Therefore the max dimension of the symmetry group is 6. If it is 6 the distribution is homogeneous.

- The theorem on the all possible Lie algebras of vanishing at 0 infinitesimal symmetries implies that b⁺_{1:1} is the only such symmetry. Therefore the max dimension of the symmetry group is 6. If it is 6 the distribution is homogeneous.
- If it is so then m = 2:

 $C(x) = \pm (x_1^2 + x_2^2)^2 + r(x_1^2 + x_2^2)^3 + higher quasi-degree$

r can be arbitrary and h.o.t. are uniquely determined by r.

The theorem on the all possible Lie algebras of vanishing at 0 infinitesimal symmetries implies that b⁺_{1:1} is the only such symmetry. Therefore the max dimension of the symmetry group is 6. If it is 6 the distribution is homogeneous.

• If it is so then
$$m = 2$$
:

 $C(x) = \pm (x_1^2 + x_2^2)^2 + r(x_1^2 + x_2^2)^3 + \text{higher quasi-degree}$

r can be arbitrary and h.o.t. are uniquely determined by r.

By a simple work with Jacobi identity the 6-dim symmetry algebra is defined by the equations

$$\begin{split} & [a_1, a_2] = a_3, \ [a_1, a_3] = a_4, \ [a_2, a_3] = a_5 \\ & [a_1, b_{1:1}^+] = a_2, \ [a_2, b_{1:1}^+] = -a_1 \\ & [a_1, a_4] = [a_2, a_5] = Ba_3 + Cb_{1:1}^+, \ [a_1, a_5] = [a_2, a_4] = 0 \end{split}$$

$$\lambda = \frac{B^2}{C}$$

is an invariant,



$$\lambda = \frac{B^2}{C}$$

is an invariant,

but unlike the case of 7-dim symmetry group it is an invariant NOT of the Lie algebra:

▲□▶ ▲□▶ ▲三▶ ▲三▶ - 三 - のへで

$$\lambda = \frac{B^2}{C}$$

is an invariant,

- but unlike the case of 7-dim symmetry group it is an invariant NOT of the Lie algebra:
- ▶ it is an invariant of the generating 2-plane span(a₁, a₂) in the 6-dim algebra.

$$\lambda = \frac{B^2}{C}$$

is an invariant,

- but unlike the case of 7-dim symmetry group it is an invariant NOT of the Lie algebra:
- ▶ it is an invariant of the generating 2-plane span(a₁, a₂) in the 6-dim algebra.

► Here: "generating 2-plane": plane span(a₁, a₂) such that a₁, a₂, [a₁, a₂], [a₁, [a₁, a₂]], [a₂, [a₁, a₂]] are linearly independent.

$$\lambda = \frac{B^2}{C}$$

is an invariant,

- but unlike the case of 7-dim symmetry group it is an invariant NOT of the Lie algebra:
- ▶ it is an invariant of the generating 2-plane span(a₁, a₂) in the 6-dim algebra.
- ► Here: "generating 2-plane": plane span(a₁, a₂) such that a₁, a₂, [a₁, a₂], [a₁, [a₁, a₂]], [a₂, [a₁, a₂]] are linearly independent.
- A distribution with transitively acting symmetry group of dimension 5 or more is uniquely determined by a generating 2-plane in the symmetry algebra.

• $C(x) = \pm (x_1^2 + x_2^2)^2 + r(x_1^2 + x_2^2)^3 + \text{higher quasi-degree}$

・ロト ・ 日 ・ モ ・ ト ・ モ ・ うへつ

• $C(x) = \pm (x_1^2 + x_2^2)^2 + r(x_1^2 + x_2^2)^3 + \text{higher quasi-degree}$

 The 6-dimensional algebra itself is one of so₃(ℝ) ⊕ so₃(ℝ) (generic value of λ = B²/C) so₃(ℝ)⊕ solvable 3-dim algebra so₃(ℝ)+ Abelian 3-dim algebra

- $C(x) = \pm (x_1^2 + x_2^2)^2 + r(x_1^2 + x_2^2)^3 + \text{higher quasi-degree}$
- The 6-dimensional algebra itself is one of so₃(ℝ) ⊕ so₃(ℝ) (generic value of λ = B²/C) so₃(ℝ)⊕ solvable 3-dim algebra so₃(ℝ)+ Abelian 3-dim algebra
- Like in the case of 7-dim symmetry group λ = B²/C is uniquely determined by r ∈ ℝ;

- $C(x) = \pm (x_1^2 + x_2^2)^2 + r(x_1^2 + x_2^2)^3 + \text{higher quasi-degree}$
- The 6-dimensional algebra itself is one of so₃(ℝ) ⊕ so₃(ℝ) (generic value of λ = B²/C) so₃(ℝ)⊕ solvable 3-dim algebra so₃(ℝ)+ Abelian 3-dim algebra
- Like in the case of 7-dim symmetry group λ = B²/C is uniquely determined by r ∈ ℝ;
- and exactly like in the case of 7-dim symmetry algebra λ takes all values except

$$\lambda = \frac{100}{9}$$

if $\lambda = 100/9$ then the distribution is flat.

- $C(x) = \pm (x_1^2 + x_2^2)^2 + r(x_1^2 + x_2^2)^3 + \text{higher quasi-degree}$
- The 6-dimensional algebra itself is one of so₃(ℝ) ⊕ so₃(ℝ) (generic value of λ = B²/C) so₃(ℝ)⊕ solvable 3-dim algebra so₃(ℝ)+ Abelian 3-dim algebra
- Like in the case of 7-dim symmetry group λ = B²/C is uniquely determined by r ∈ ℝ;
- and exactly like in the case of 7-dim symmetry algebra λ takes all values except

$$\lambda = \frac{100}{9}$$

if $\lambda = 100/9$ then the distribution is flat.

► Cartan has these results in Ch. XI. For the case so₃(ℝ) ⊕ so₃(ℝ) he has a normal form with parameters m, n ∈ ℝ whose ratio m/n is an invariant, with exceptional value 81. • Why $\frac{100}{9}$ for both 7-dim and 6-dim algebra?

・ロト・(型ト・(型ト・(型ト))
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・</li

- Why $\frac{100}{9}$ for both 7-dim and 6-dim algebra?
- Some mystery, but exactly the same as the mystery in the problem of a ball of radius R₁ rolling over another ball of radius R₂ without slipping or twisting. This rolling is described by (2, 3, 5) distribution with clear (natural) 6-dim group of symmetries so₃(ℝ) ⊕ so₃(ℝ). One of them is b⁺_{1:1} up to a diffeomorphism. The presence of b⁺_{1:1} implies that either there are no other symmetries or the distribution is flat, i.e. there are 8 other symmetries. The distribution is a generating plane in so₃(ℝ) ⊕ so₃(ℝ) and its invariant is

$$\lambda = B^2/C = \left(\frac{\mu^2 + 1}{\mu}\right)^2, \ \mu = \frac{R_1}{R_2}$$

- Why $\frac{100}{9}$ for both 7-dim and 6-dim algebra?
- Some mystery, but exactly the same as the mystery in the problem of a ball of radius R₁ rolling over another ball of radius R₂ without slipping or twisting. This rolling is described by (2, 3, 5) distribution with clear (natural) 6-dim group of symmetries so₃(ℝ) ⊕ so₃(ℝ). One of them is b⁺_{1:1} up to a diffeomorphism. The presence of b⁺_{1:1} implies that either there are no other symmetries or the distribution is flat, i.e. there are 8 other symmetries. The distribution is a generating plane in so₃(ℝ) ⊕ so₃(ℝ) and its invariant is

$$\lambda = B^2/C = \left(\frac{\mu^2 + 1}{\mu}\right)^2, \ \mu = \frac{R_1}{R_2}$$

• The equation $\left(\frac{\mu^2+1}{\mu}\right)^2 = \frac{100}{9}$ has two positive solutions $\mu = 3$ and 1/3 which gives one more proof of

うしん 山 くはゃくはゃくむゃくしゃ

Theorem (R. Bryant). The distribution is flat if and only if the ratio of the radii is 3 (or 1/3).

Recent works where this theorem was proved $+\ very$ interesting related math is explained:

A.Agrachev, Rolling Balls and Octonions, 2007

G.Bor, R. Montgomery, G₂ and the rolling distribution, 2009

▲□▶ ▲□▶ ▲ 臣▶ ★ 臣▶ 三臣 - のへぐ

wrt weights $\lambda_1 = \lambda_2 = 1$, $\lambda_3 = 2$, $\lambda_4 = \lambda_5 = 3$

• Quasi-homogeneity wrt weights $\lambda_1 = \lambda_2 = 1$, $\lambda_3 = 2$, $\lambda_4 = \lambda_5 = 3$

Quasi-homogeneous degree i vector field = linear combination of quasi-degree i monomial vector fields.

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● ● ●

- Quasi-homogeneous degree i vector field = linear combination of quasi-degree i monomial vector fields.
- Quasi-degree of a monomial vector field $x^{\alpha} \frac{\partial}{\partial x_i}$, is $(\lambda, \alpha) \lambda_j$

・ロト ・ 同 ト ・ 三 ト ・ 三 ・ うへつ

wrt weights $\lambda_1 = \lambda_2 = 1$, $\lambda_3 = 2$, $\lambda_4 = \lambda_5 = 3$

- Quasi-homogeneous degree i vector field = linear combination of quasi-degree i monomial vector fields.
- Quasi-degree of a monomial vector field $x^{\alpha} \frac{\partial}{\partial x_i}$, is $(\lambda, \alpha) \lambda_j$

・ロト ・ 同 ト ・ 三 ト ・ 三 ・ うへつ

► The Lie brackets respect quasi-homogeneity: [(i), (j)] ∈ (i + j).

wrt weights $\lambda_1 = \lambda_2 = 1$, $\lambda_3 = 2$, $\lambda_4 = \lambda_5 = 3$

- Quasi-homogeneous degree i vector field = linear combination of quasi-degree i monomial vector fields.
- Quasi-degree of a monomial vector field $x^{\alpha} \frac{\partial}{\partial x_i}$, is $(\lambda, \alpha) \lambda_j$
- ► The Lie brackets respect quasi-homogeneity: [(i), (j)] ∈ (i + j).
- ► Except the zero vector field there are no quasi-homogeneous degree ≤ -4 vector fields. The quasi-homogeneous vector fields of degree -3 have the form a ∂/∂x₃ + b ∂/∂x₄

・ロト ・ 同 ト ・ 三 ト ・ 三 ・ うへつ

wrt weights $\lambda_1 = \lambda_2 = 1$, $\lambda_3 = 2$, $\lambda_4 = \lambda_5 = 3$

- Quasi-homogeneous degree i vector field = linear combination of quasi-degree i monomial vector fields.
- Quasi-degree of a monomial vector field $x^{\alpha} \frac{\partial}{\partial x_i}$, is $(\lambda, \alpha) \lambda_j$
- ► The Lie brackets respect quasi-homogeneity: [(i), (j)] ∈ (i + j).
- ► Except the zero vector field there are no quasi-homogeneous degree ≤ -4 vector fields. The quasi-homogeneous vector fields of degree -3 have the form a ∂/∂x₃ + b ∂/∂x₄
- ► It follows that any length ≥ 4 bracket of quasi-homogeneous degree -1 vector fields is 0.

wrt weights $\lambda_1 = \lambda_2 = 1$, $\lambda_3 = 2$, $\lambda_4 = \lambda_5 = 3$

- Quasi-homogeneous degree i vector field = linear combination of quasi-degree i monomial vector fields.
- Quasi-degree of a monomial vector field $x^{\alpha} \frac{\partial}{\partial x_i}$, is $(\lambda, \alpha) \lambda_j$
- ► The Lie brackets respect quasi-homogeneity: [(i), (j)] ∈ (i + j).
- ► Except the zero vector field there are no quasi-homogeneous degree ≤ -4 vector fields. The quasi-homogeneous vector fields of degree -3 have the form a ∂/∂x₃ + b ∂/∂x₄
- ► It follows that any length ≥ 4 bracket of quasi-homogeneous degree -1 vector fields is 0.

► Consequently, any (2,3,5) distribution spanned by quasi-homogeneous degree -1 vector fields is flat. Nilpotent approximation (starting point for the normalization)

・ロト ・ 日 ・ モ ト ・ モ ・ うへぐ

(starting point for the normalization)

There are many ways to prove the following statement (particular case of much more general Bellaiche theorem for arbitrary bracket generating tuples of vector fields):

Proposition. Fix any (2,3,5) distribution spanned by quasi-homogeneous degree -1 vector fields $\mathcal{N}_1, \mathcal{N}_2$. Any (2,3,5) distribution D is diffeomorphic to a distribution spanned by vector fields of the form

$$V_1 = \mathcal{N}_1 + \cdots, \ V_2 = \mathcal{N}_2 + \cdots,$$

where the dots denote terms of quasi-degree ≥ 0 .

(starting point for the normalization)

There are many ways to prove the following statement (particular case of much more general Bellaiche theorem for arbitrary bracket generating tuples of vector fields):

Proposition. Fix any (2,3,5) distribution spanned by quasi-homogeneous degree -1 vector fields $\mathcal{N}_1, \mathcal{N}_2$. Any (2,3,5) distribution D is diffeomorphic to a distribution spanned by vector fields of the form

$$V_1 = \mathcal{N}_1 + \cdots, \ V_2 = \mathcal{N}_2 + \cdots,$$

where the dots denote terms of quasi-degree ≥ 0 .

(starting point for the normalization)

There are many ways to prove the following statement (particular case of much more general Bellaiche theorem for arbitrary bracket generating tuples of vector fields):

Proposition. Fix any (2,3,5) distribution spanned by quasi-homogeneous degree -1 vector fields $\mathcal{N}_1, \mathcal{N}_2$. Any (2,3,5) distribution D is diffeomorphic to a distribution spanned by vector fields of the form

$$V_1 = \mathcal{N}_1 + \cdots, \ V_2 = \mathcal{N}_2 + \cdots,$$

where the dots denote terms of quasi-degree ≥ 0 .

► The distribution span(N₁, N₂) (or the couple N = (N₁, N₂)) is the nilpotent approximation of D.

(starting point for the normalization)

There are many ways to prove the following statement (particular case of much more general Bellaiche theorem for arbitrary bracket generating tuples of vector fields):

Proposition. Fix any (2,3,5) distribution spanned by quasi-homogeneous degree -1 vector fields $\mathcal{N}_1, \mathcal{N}_2$. Any (2,3,5) distribution D is diffeomorphic to a distribution spanned by vector fields of the form

$$V_1 = \mathcal{N}_1 + \cdots, \ V_2 = \mathcal{N}_2 + \cdots,$$

where the dots denote terms of quasi-degree ≥ 0 .

- ► The distribution span(N₁, N₂) (or the couple N = (N₁, N₂)) is the nilpotent approximation of D.
- It is the symbol of D (graded nilpotent (2,3,5) Lie algebra; all such symbols are isomorphic) represented by vector fields.

Infinitesimal linear operator

associated with the nilpotent approximation $\mathcal{N} = (\mathcal{N}_1, \mathcal{N}_2)$.

▲□▶ ▲□▶ ▲□▶ ▲□▶ □ ● ● ●

Infinitesimal linear operator

►

associated with the nilpotent approximation $\mathcal{N} = (\mathcal{N}_1, \mathcal{N}_2)$.

- A 2-distribution can be treated as a couple of vector fields V = (V₁, V₂) and two distributions are diffeomorphic if the corresponding couples of vector fields can be brought one to the other by a local diffeomorphism Φ and multiplication by a non-singular 2 × 2 matrix H = H(x).
- ▶ The infinitesimal linear operator associated with \mathcal{N} (in what follows infinitesimal operator; notation $L_{\mathcal{N}}$ is the linearization (at the identity transformation) of the map $(\Phi, H) \rightarrow H\Phi * \mathcal{N}$. It is a map from the Lie algebra of the (pseudo)-group $\{\Phi, H\}$ which is (Z, h), where Z is a vector field h = h(x) is any 2×2 matrix, to the space of couples of vector fields.

Infinitesimal linear operator

►

associated with the nilpotent approximation $\mathcal{N} = (\mathcal{N}_1, \mathcal{N}_2)$.

- A 2-distribution can be treated as a couple of vector fields V = (V₁, V₂) and two distributions are diffeomorphic if the corresponding couples of vector fields can be brought one to the other by a local diffeomorphism Φ and multiplication by a non-singular 2 × 2 matrix H = H(x).
- ▶ The infinitesimal linear operator associated with \mathcal{N} (in what follows infinitesimal operator; notation $L_{\mathcal{N}}$ is the linearization (at the identity transformation) of the map $(\Phi, H) \rightarrow H\Phi * \mathcal{N}$. It is a map from the Lie algebra of the (pseudo)-group $\{\Phi, H\}$ which is (Z, h), where Z is a vector field h = h(x) is any 2×2 matrix, to the space of couples of vector fields.

$$L_{\mathcal{N}}: (Z,h) \rightarrow [Z,\mathcal{N}] + h\mathcal{N}.$$

► The operator L_N respects the quasi-homogeneous filtration: if h and Z are quasi-homogeneous of degree i then L_N(Z, h) is a couple of quasi-homogeneous vector fields of degree i - 1.

◆□▶ ◆□▶ ◆ □▶ ◆ □▶ - □ - のへぐ

- ► The operator L_N respects the quasi-homogeneous filtration: if h and Z are quasi-homogeneous of degree i then L_N(Z, h) is a couple of quasi-homogeneous vector fields of degree i - 1.
- The vector field part of

$$kerL_{\mathcal{N}} = kerL_{\mathcal{N}}^{(-3)} + kerL_{\mathcal{N}}^{(-2)} + \cdots$$

is the Lie algebra of infinitesimal symmetries of \mathcal{N} . The matrix part of $kerL_{\mathcal{N}}$ is uniquely determined by the vector field part.

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● ● ●

- ► The operator L_N respects the quasi-homogeneous filtration: if h and Z are quasi-homogeneous of degree i then L_N(Z, h) is a couple of quasi-homogeneous vector fields of degree i - 1.
- The vector field part of

$$kerL_{\mathcal{N}} = kerL_{\mathcal{N}}^{(-3)} + kerL_{\mathcal{N}}^{(-2)} + \cdots$$

is the Lie algebra of infinitesimal symmetries of \mathcal{N} . The matrix part of $kerL_{\mathcal{N}}$ is uniquely determined by the vector field part. **Dimension of** $kerL_{\mathcal{N}}^{(i)}$:

i ≥ 4: 0

▶ i = -1: 2 i = -2: 1 i = -3: 2 i = 0: 4 i = 1: 2 i = 1: 2 i = 2: 1 i = 3: 2

- \blacktriangleright The operator L_N respects the quasi-homogeneous filtration: if h and Z are quasi-homogeneous of degree i then $L_{\mathcal{N}}(Z, h)$ is a couple of quasi-homogeneous vector fields of degree i - 1.
- The vector field part of

$$kerL_{\mathcal{N}} = kerL_{\mathcal{N}}^{(-3)} + kerL_{\mathcal{N}}^{(-2)} + \cdots$$

is the Lie algebra of infinitesimal symmetries of \mathcal{N} . The matrix part of $kerL_N$ is uniquely determined by the vector field part. **Dimension of** $kerL_{N}^{(i)}$:

i > 4: 0

i = -1: 2i = 1: 2i = -2: 1 i = 0: 4i = 2: 1 i = -3: 2 i = 3: 2

Type of vector fields

Non-vanishing quasi-linear with zero quasi-linear part

< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > <

- ► The operator L_N respects the quasi-homogeneous filtration: if h and Z are quasi-homogeneous of degree i then L_N(Z, h) is a couple of quasi-homogeneous vector fields of degree i - 1.
- The vector field part of

$$kerL_{\mathcal{N}} = kerL_{\mathcal{N}}^{(-3)} + kerL_{\mathcal{N}}^{(-2)} + \cdots$$

is the Lie algebra of infinitesimal symmetries of \mathcal{N} . The matrix part of $kerL_{\mathcal{N}}$ is uniquely determined by the vector field part.

Dimension of $kerL_{\mathcal{N}}^{(i)}$:

i ≥ 4: 0

i = -1: 2i = -2: 1i = -3: 2i = 0: 4i = 1: 2i = 1: 2i = 2: 1i = 3: 2i = 3: 2i = 1: 2i = 2: 1i = 3: 2i = 3: 3i = 3

Type of vector fields

Non-vanishing quasi-linear with zero quasi-linear part
 Lie brackets: [(i), (j)] ∈ (i + j), |i| ≥ 4 ⇒ (i) = {0}

► Lie algebra
$$kerL_{N} = \sum_{i=-3}^{3} kerL_{N}^{(i)}$$
:
 $i = -1: dim = 2$ $i = 0: dim = 4$ $i = 1: dim = 2$
 $i = -2: dim = 1$ $i = 2: dim = 3$
 $i = -3: dim = 2$ $i = 3: dim = 1$

(2,3,5) nilpotent + $gl_2(\mathbb{R})$ + (2,3,5) nilpotent

► Lie algebra
$$kerL_{\mathcal{N}} = \sum_{i=-3}^{3} kerL_{\mathcal{N}}^{(i)}$$
:
 $i = -1: dim = 2$ $i = 0: dim = 4$ $i = 1: dim = 2$
 $i = -2: dim = 1$ $i = 2: dim = 3$
 $i = -3: dim = 2$ $i = 3: dim = 1$

(2,3,5) nilpotent + $gl_2(\mathbb{R})$ + (2,3,5) nilpotent

► It is the 14-dim algebra g₂ whose negative part ((i = -1, -2, -3) consists of non-vanishing symmetries, central part (i = 0) consists of quasi-linear infinitesimal symmetries, and the positive part (i = 1, 2, 3) consists of vanishing at 0 symmetries with zero quasi-linear approximation.

► Lie algebra
$$kerL_{N} = \sum_{i=-3}^{3} kerL_{N}^{(i)}$$
:
 $i = -1: dim = 2$ $i = 0: dim = 4$ $i = 1: dim = 2$
 $i = -2: dim = 1$ $i = 2: dim = 3$
 $i = -3: dim = 2$ $i = 3: dim = 1$

(2,3,5) nilpotent + $gl_2(\mathbb{R})$ + (2,3,5) nilpotent

- ► It is the 14-dim algebra g₂ whose negative part ((i = -1, -2, -3) consists of non-vanishing symmetries, central part (i = 0) consists of quasi-linear infinitesimal symmetries, and the positive part (i = 1, 2, 3) consists of vanishing at 0 symmetries with zero quasi-linear approximation.
- The last two pages are "isomorphic" to g₂ in terms of Tanaka prolongation.

W⁽ⁱ⁾ = a complementary space to the image of L⁽ⁱ⁾_N (of our choice).

- W⁽ⁱ⁾ = a complementary space to the image of L⁽ⁱ⁾_N (of our choice).
- Like in all local classification problems, the first normalization step gives the preliminary normal form

$$V = \mathcal{N} + V^{(0)} + V^{(1)} + \cdots,$$

 $V^{(i)} = (V_1, V_2) \in W^{(i)}$

serving for all (2,3,5) distributions with a fixed nilpotent approximation $\ensuremath{\mathcal{N}}.$

It is easy to compute that

$$W^{(0)} = W^{(1)} = W^{(2)} = \{0\}, \quad W^{(3)} \neq 0.$$

W⁽ⁱ⁾ = a complementary space to the image of L⁽ⁱ⁾_N (of our choice).

- W⁽ⁱ⁾ = a complementary space to the image of L⁽ⁱ⁾_N (of our choice).
- ▶ Therefore any (2,3,5) distribution is diffeomorphic to

$$V = \mathcal{N} + V^{(d \ge 3)} + V^{(d+1)} + \cdots,$$

$$V^{(i)} = (V_1, V_2) \in W^{(i)}, \quad V^{(d)} \neq 0.$$

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□

or to its nilpotent approximation \mathcal{N} .

- W⁽ⁱ⁾ = a complementary space to the image of L⁽ⁱ⁾_N (of our choice).
- ▶ Therefore any (2,3,5) distribution is diffeomorphic to

$$V = \mathcal{N} + V^{(d \ge 3)} + V^{(d+1)} + \cdots,$$

$$V^{(i)} = (V_1, V_2) \in W^{(i)}, \quad V^{(d)} \neq 0.$$

▶ Proposition. An analytic distribution is flat if and only if it is formally diffeomorphic to N.

・ロト・日本・モート モー うへぐ

- W⁽ⁱ⁾ = a complementary space to the image of L⁽ⁱ⁾_N (of our choice).
- Therefore any (2,3,5) distribution is diffeomorphic to

$$V = \mathcal{N} + V^{(d \ge 3)} + V^{(d+1)} + \cdots,$$

$$V^{(i)} = (V_1, V_2) \in W^{(i)}, \quad V^{(d)} \neq 0.$$

- ▶ Proposition. An analytic distribution is flat if and only if it is formally diffeomorphic to *N*.
- ▶ Proposition. d ≥ 3 is an invariant of a non-flat distribution (I call it degree of non-flatness).

- W⁽ⁱ⁾ = a complementary space to the image of L⁽ⁱ⁾_N (of our choice).
- ▶ Therefore any (2,3,5) distribution is diffeomorphic to

$$V = \mathcal{N} + V^{(d \ge 3)} + V^{(d+1)} + \cdots,$$

$$V^{(i)} = (V_1, V_2) \in W^{(i)}, \quad V^{(d)} \neq 0.$$

- ▶ Proposition. An analytic distribution is flat if and only if it is formally diffeomorphic to *N*.
- ▶ Proposition. d ≥ 3 is an invariant of a non-flat distribution (I call it degree of non-flatness).
- ▶ Proposition. V^(d≥3) modulo action of quasi-linear symmetries of N is invariant. I call generalized Cartan invariant.

- W⁽ⁱ⁾ = a complementary space to the image of L⁽ⁱ⁾_N (of our choice).
- Therefore any (2,3,5) distribution is diffeomorphic to

$$V = \mathcal{N} + V^{(d \ge 3)} + V^{(d+1)} + \cdots,$$

$$V^{(i)} = (V_1, V_2) \in W^{(i)}, \quad V^{(d)} \neq 0.$$

- ▶ Proposition. An analytic distribution is flat if and only if it is formally diffeomorphic to N.
- ▶ Proposition. d ≥ 3 is an invariant of a non-flat distribution (I call it degree of non-flatness).
- ▶ Proposition. V^(d≥3) modulo action of quasi-linear symmetries of N is invariant. I call generalized Cartan invariant.

Claim. If d = 3 then the generalized Cartan invariant is exactly the Cartan invariant at 0 ∈ ℝ⁵. W⁽ⁱ⁾ = a complementary space to the image of L⁽ⁱ⁾_N (of our choice).

・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・
 ・

- W⁽ⁱ⁾ = a complementary space to the image of L⁽ⁱ⁾_N (of our choice).
- ► Finding *W* reduces to finding a complementary space to the image of the linear operator

$$T: (f_1, f_2, g) \rightarrow \begin{pmatrix} \mathcal{N}_1(f_1) + g & \mathcal{N}_1(f_2) \\ \mathcal{N}_2(f_1) & \mathcal{N}_2(f_2) + g \end{pmatrix}$$

◆□▶ ◆□▶ ◆ □▶ ◆ □▶ - □ - のへぐ

where f_1, f_2, g are function germs:

- W⁽ⁱ⁾ = a complementary space to the image of L⁽ⁱ⁾_N (of our choice).
- ► Finding *W* reduces to finding a complementary space to the image of the linear operator

$$T: (f_1, f_2, g) \rightarrow \begin{pmatrix} \mathcal{N}_1(f_1) + g & \mathcal{N}_1(f_2) \\ \mathcal{N}_2(f_1) & \mathcal{N}_2(f_2) + g \end{pmatrix}$$

where f_1, f_2, g are function germs:

► if U is a complementary space to the image of T in the space Mat_{2×2}(x) then

$$\left\{ A(x) \begin{pmatrix} [N_1, [\mathcal{N}_1, \mathcal{N}_2]] \\ [N_2, [\mathcal{N}_1, \mathcal{N}_2]] \end{pmatrix}, \quad A(x) \in U \right\}$$

is a complementary space to the image of L_N .

• Choice of nilpotent approximation $\mathcal{N} = span(\mathcal{N}_1, \mathcal{N}_2)$:

$$\mathcal{N}_{1} = \frac{\partial}{\partial x_{1}} + x_{2} \left(\frac{\partial}{\partial x_{3}} + x_{1} \frac{\partial}{\partial x_{4}} + x_{2} \frac{\partial}{\partial x_{5}} \right)$$
$$\mathcal{N}_{2} = \frac{\partial}{\partial x_{2}} - x_{1} \left(\frac{\partial}{\partial x_{3}} + x_{1} \frac{\partial}{\partial x_{4}} + x_{2} \frac{\partial}{\partial x_{5}} \right)$$

・ロト ・ 日 ・ モ ト ・ モ ・ うへぐ

• Choice of nilpotent approximation $\mathcal{N} = span(\mathcal{N}_1, \mathcal{N}_2)$:

$$\mathcal{N}_{1} = \frac{\partial}{\partial x_{1}} + x_{2} \left(\frac{\partial}{\partial x_{3}} + x_{1} \frac{\partial}{\partial x_{4}} + x_{2} \frac{\partial}{\partial x_{5}} \right)$$
$$\mathcal{N}_{2} = \frac{\partial}{\partial x_{2}} - x_{1} \left(\frac{\partial}{\partial x_{3}} + x_{1} \frac{\partial}{\partial x_{4}} + x_{2} \frac{\partial}{\partial x_{5}} \right)$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

► Advantage: Any quasi-linear symmetry of span(N₁, N₂) (= symmetry of quasi-degree 0) is linear. • Choice of nilpotent approximation $\mathcal{N} = span(\mathcal{N}_1, \mathcal{N}_2)$:

$$\mathcal{N}_{1} = \frac{\partial}{\partial x_{1}} + x_{2} \left(\frac{\partial}{\partial x_{3}} + x_{1} \frac{\partial}{\partial x_{4}} + x_{2} \frac{\partial}{\partial x_{5}} \right)$$
$$\mathcal{N}_{2} = \frac{\partial}{\partial x_{2}} - x_{1} \left(\frac{\partial}{\partial x_{3}} + x_{1} \frac{\partial}{\partial x_{4}} + x_{2} \frac{\partial}{\partial x_{5}} \right)$$

- ► Advantage: Any quasi-linear symmetry of span(N₁, N₂) (= symmetry of quasi-degree 0) is linear.
- It has the form

$$g_Q: egin{pmatrix} x_1 \ x_2 \end{pmatrix} o Q egin{pmatrix} x_1 \ x_2 \end{pmatrix}, \ x_3 o det Q \cdot x_3, \ egin{pmatrix} x_4 \ x_5 \end{pmatrix} o det Q \cdot Q egin{pmatrix} x_4 \ x_5 \end{pmatrix}$$

- ロ ト - 4 回 ト - 4 □

where Q is a non-singular 2 \times 2 matrix.

Asymptotically exact normal form

$$V = \mathcal{N} + C(x) \begin{pmatrix} x_1 x_2 & x_2^2 \\ -x_1^2 & -x_1 x_2 \end{pmatrix} \begin{pmatrix} \frac{\partial}{\partial x_4} \\ \frac{\partial}{\partial x_5} \end{pmatrix}$$

・ロト ・ 日 ・ モ ト ・ モ ・ うへぐ

Asymptotically exact normal form

$$V = \mathcal{N} + C(x) \begin{pmatrix} x_1 x_2 & x_2^2 \\ -x_1^2 & -x_1 x_2 \end{pmatrix} \begin{pmatrix} \frac{\partial}{\partial x_4} \\ \frac{\partial}{\partial x_5} \end{pmatrix}$$

► Almost exact normal form: same with C(x) in the ideal generated by the monomials

Monomials generating the ideal I	Quasi-degree
$x_1^i x_2^j, \ i+j=4$	4
$x_1^i x_2^j x_3, \ i+j=3$	5
$x_1^i x_2^j x_3^2, i+j=2$	6
$x_1x_3\theta, x_2x_3\theta$	7
θ^2	8

$$\theta = x_1 x_4 - x_2 x_5$$

▲□▶ ▲□▶ ▲三▶ ▲三▶ - 三 - のへで



$$V=\mathcal{N}+\mathcal{C}(x)\begin{pmatrix} x_1x_2 & x_2^2\\ -x_1^2 & -x_1x_2 \end{pmatrix}\begin{pmatrix} \frac{\partial}{\partial x_4}\\ \frac{\partial}{\partial x_5} \end{pmatrix}, \ \ \mathcal{C}(x)\in \mathsf{the ideal}.$$

▲□▶ ▲□▶ ▲三▶ ▲三▶ 三三 - のへで



$$V = \mathcal{N} + C(x) \begin{pmatrix} x_1 x_2 & x_2^2 \\ -x_1^2 & -x_1 x_2 \end{pmatrix} \begin{pmatrix} \frac{\partial}{\partial x_4} \\ \frac{\partial}{\partial x_5} \end{pmatrix}, \quad C(x) \in \mathsf{the ideal}.$$

• The expansion of C(x) by quasi-homogeneous terms:

$$C(x) = C^{(d+1)}(x) + C^{(d+2)}(x) + \cdots, \ d \ge 3$$

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□

where d is the degree of non-flatness and $C^{d+1}(x)$ is the generalized Cartan invariant.



$$V = \mathcal{N} + C(x) \begin{pmatrix} x_1 x_2 & x_2^2 \\ -x_1^2 & -x_1 x_2 \end{pmatrix} \begin{pmatrix} \frac{\partial}{\partial x_4} \\ \frac{\partial}{\partial x_5} \end{pmatrix}, \quad C(x) \in \mathsf{the ideal}.$$

• The expansion of C(x) by quasi-homogeneous terms:

$$C(x) = C^{(d+1)}(x) + C^{(d+2)}(x) + \cdots, \ d \ge 3$$

where d is the degree of non-flatness and $C^{d+1}(x)$ is the generalized Cartan invariant.

• A quasi-linear symmetry g_Q of \mathcal{N} brings C(x) to $C(g_Q(x))$.

- ロ ト - 4 回 ト - 4 □



$$V = \mathcal{N} + C(x) \begin{pmatrix} x_1 x_2 & x_2^2 \\ -x_1^2 & -x_1 x_2 \end{pmatrix} \begin{pmatrix} \frac{\partial}{\partial x_4} \\ \frac{\partial}{\partial x_5} \end{pmatrix}, \quad C(x) \in \mathsf{the ideal}.$$

• The expansion of C(x) by quasi-homogeneous terms:

$$C(x) = C^{(d+1)}(x) + C^{(d+2)}(x) + \cdots, \ d \ge 3$$

where d is the degree of non-flatness and $C^{d+1}(x)$ is the generalized Cartan invariant.

- A quasi-linear symmetry g_Q of \mathcal{N} brings C(x) to $C(g_Q(x))$.
- ► C⁽⁴⁾(x) = C⁽⁴⁾(x₁, x₂), therefore if d = 3 (minimal possible degree of non-flatness) then the generalized invariant is the classical Cartan invariant.

If we replace $0 \in \mathbb{R}^5$ by as point nearby, the Cartan tensor (defined in this way) changes in very involved way, but its infinitesimal change is simple. This allows to give a simple proof of Cartan theorem that the distribution is flat if and only if the Cartan tensor vanishes identically reducing it to Frobenius theorem.

► If b is an infinitesimal symmetry vanishing at 0 then its quasi-linear part V_A annihilates the generalized Cartan invariant C^(d+1)(x). Since C^(d+1)(x) is a non-zero polynomial, it is possible only if V_A is a resonant vector field: either one of the eigenvalues is 0 or their ratio is a negative integer number. This is the main point for all theorems on infinitesimal symmetries vanishing at 0.

- ► If b is an infinitesimal symmetry vanishing at 0 then its quasi-linear part V_A annihilates the generalized Cartan invariant C^(d+1)(x). Since C^(d+1)(x) is a non-zero polynomial, it is possible only if V_A is a resonant vector field: either one of the eigenvalues is 0 or their ratio is a negative integer number. This is the main point for all theorems on infinitesimal symmetries vanishing at 0.
- ▶ The proof of linearization theorems for such symmetries follows from the second normalization step in which the almost exact normal is "corrected" using the 5-dimensional "positive" part of g₂.

Thanks again to the Organizers for their hard work for us before and during the conference

・ロト ・ 日 ・ モ ト ・ モ ・ うへぐ

Thanks again to the Organizers for their hard work for us before and during the conference

◆□▶ ◆□▶ ◆ □▶ ◆ □▶ - □ - のへぐ

 Many more years of fruitful math to Professor Keizo Yamaguchi and Professor Reiko Miyaoka

- Thanks again to the Organizers for their hard work for us before and during the conference
- Many more years of fruitful math to Professor Keizo Yamaguchi and Professor Reiko Miyaoka
- I am very thankful to Professor Ishikawa and Professor Morimoto for inviting me to Japan in 2005 where I got involved to this topic

- Thanks again to the Organizers for their hard work for us before and during the conference
- Many more years of fruitful math to Professor Keizo Yamaguchi and Professor Reiko Miyaoka
- I am very thankful to Professor Ishikawa and Professor Morimoto for inviting me to Japan in 2005 where I got involved to this topic
- Robert Bryant, Andrei Agracjev and especially Igor Zelenko helped me to understand much of the math I need to do this work

- Thanks again to the Organizers for their hard work for us before and during the conference
- Many more years of fruitful math to Professor Keizo Yamaguchi and Professor Reiko Miyaoka
- I am very thankful to Professor Ishikawa and Professor Morimoto for inviting me to Japan in 2005 where I got involved to this topic
- Robert Bryant, Andrei Agracjev and especially Igor Zelenko helped me to understand much of the math I need to do this work
- And thanks to the AUDIENCE!

If you had no time to read all what is written in some pages and are interested to return to them, this file will be published soon in my homepage (google Technion, Math).